

Compressible Flow - TME085

Chapter 12

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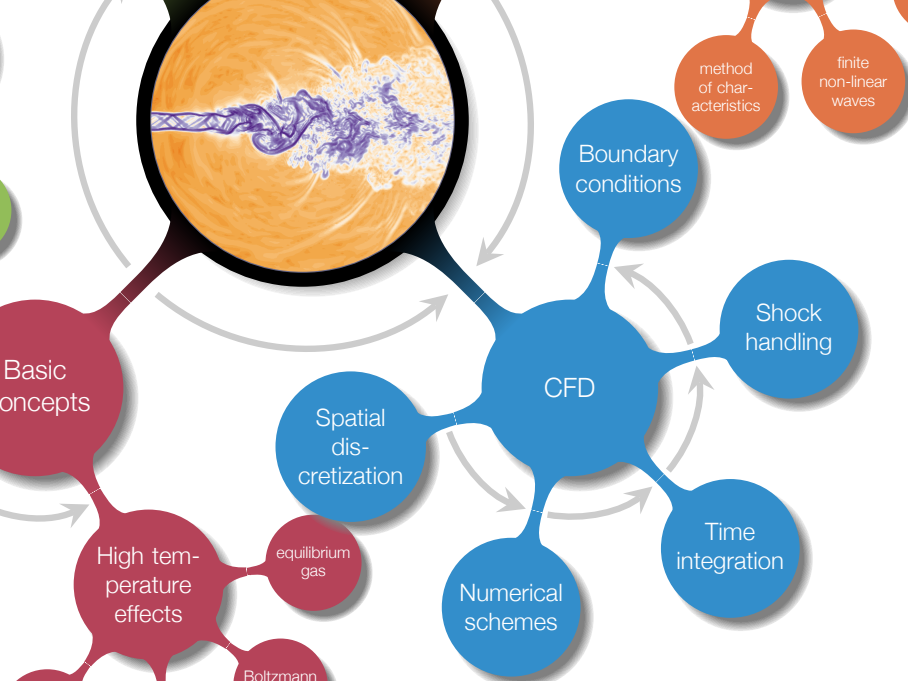
```

4 #undef __FUNCT__
5 #define __FUNCT__ "RungeKutta::fwd"
6 PetscErrorCode RungeKutta::fwd(Domain *dom){
7     PetscErrorCode ierr=0;
8
9     ierr=G3DCopy(dom->cons,cons0);CHKERRQ(ierr);
10
11     /* RK1 */
12
13     dom->update();
14
15     dcons->evaluate(dom);
16
17     ierr=G3DWAXPY(dom->cons,1.0,dcons,cons0);CHKERRQ(ierr);
18     ierr=G3DAXPBY(cons0,0.5,0.5,dom->cons);CHKERRQ(ierr);
19
20     /* RK2 */
21
22     dom->update();
23     dcons->evaluate(dom);
24
25     ierr=G3DWAXPY(dom->cons,0.5,dcons,cons0);CHKERRQ(ierr);
26
27     /* RK3 */

```

Chapter 12 - The Time-Marching Technique

Overview

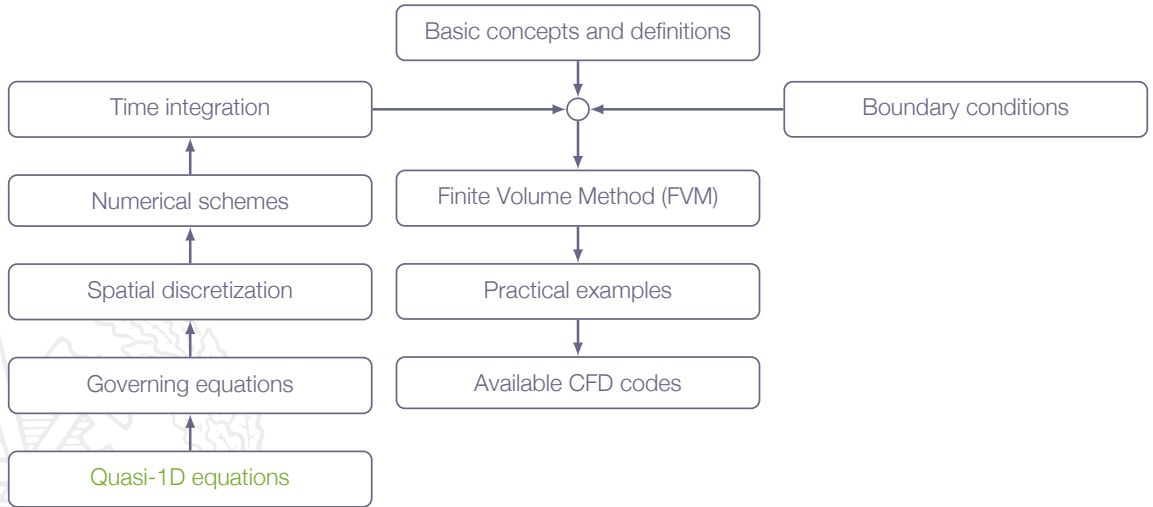


Learning Outcomes

- 12 **Explain** the main principles behind a modern Finite Volume CFD code and such concepts as explicit/implicit time stepping, CFL number, conservation, handling of compression shocks, and boundary conditions
- 14 **Analyze** and **verify** the quality of the numerical solution
- 15 **Explain** the limitations in fluid flow simulation software

time for CFD!

Roadmap - The Time-Marching Technique



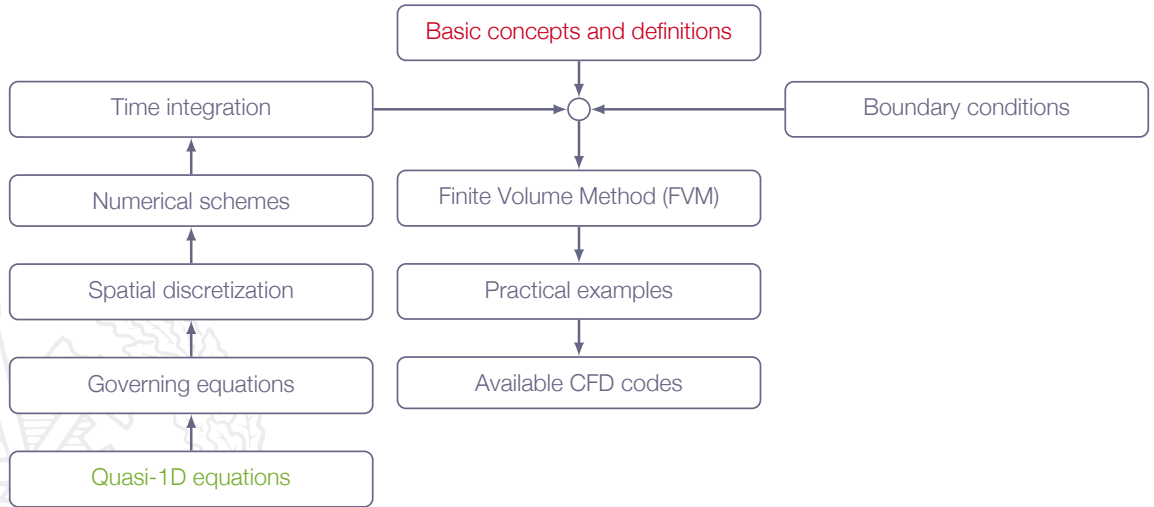
Motivation

Computational Fluid Dynamics (CFD) is the backbone of all practical engineering compressible flow analysis

As an engineer doing numerical compressible flow analyzes it is extremely important to have knowledge about the fundamental numerical principles and their **limitations**

Going through the material covered in this section will not make you understand all the details but you will get a feeling, which is a good start

Roadmap - The Time-Marching Technique



The Time-Marching Technique

The problems that we like to investigate numerically within the field of compressible flows can be categorized as

**steady-state
compressible flows**

**unsteady
compressible flows**

The **Time-marching technique** is a solver framework that addresses both problem categories

The Time-Marching Technique

Steady-state problems:

1. define simple initial solution
2. apply specified boundary conditions
3. march in time until steady-state solution is reached

Unsteady problems:

1. apply specified initial solution
2. apply specified boundary conditions
3. march in time for specified total time to reach a desired unsteady solution

establish fully developed flow before initiating data sampling

The Time-Marching Technique

The time-marching approach is a good alternative for simulating flows where there are both supersonic and subsonic regions

supersonic/hyperbolic:

- perturbations propagate in preferred directions

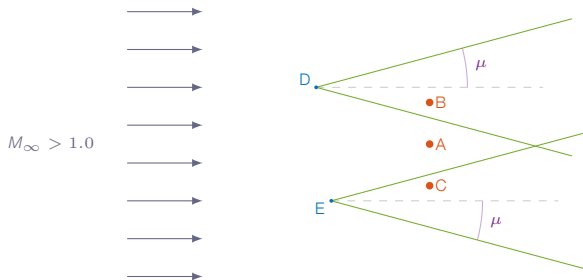
- zone of influence/zone of dependence

- PDEs can be transformed into ODEs

subsonic/elliptic:

- perturbations propagate in all directions

Zone of Influence and Zone of Dependence



A, B and C at the same axial position in the flow

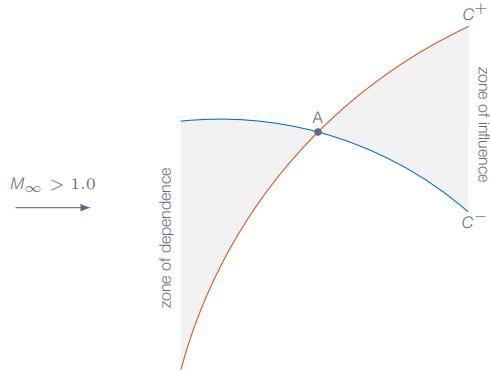
D and E are located upstream of A, B and C

Mach waves generated at D will affect the flow in B but not in A and C

Mach waves generated at E will affect the flow in C but not in A and B

The flow in A is unaffected by the both D and E

Zone of Influence and Zone of Dependence



The **zone of dependence** for point A and the **zone of influence** of point A are defined by C^+ and C^- characteristic lines

Characterization of CFD Methods

Density-based

Pressure-based

Fully coupled

Segregated

Structured

Unstructured

Explicit

Implicit

Characterization of CFD Methods

Approach taken in this presentation

Density-based

Fully coupled

Structured

Explicit

Pressure-based

Segregated

Unstructured

Implicit

Characterization of CFD Methods - Equations

Density-based

solve for density in the continuity equation

suitable for transonic/supersonic flows

Pressure-based

the continuity and momentum equations are combined to form a pressure correction equation

suitable for subsonic/transonic flows

Characterization of CFD Methods - Solver Approach

Fully coupled

all equations (continuity, momentum, energy, ...) are solved simultaneously
suitable for transonic/supersonic flows

Segregated

the governing equations are solved in sequence
suitable for subsonic flows

Characterization of CFD Methods - Time Stepping

Explicit

- short time steps
- + very stable

Implicit

- + longer time steps possible

Characterization of CFD Methods - Time Stepping

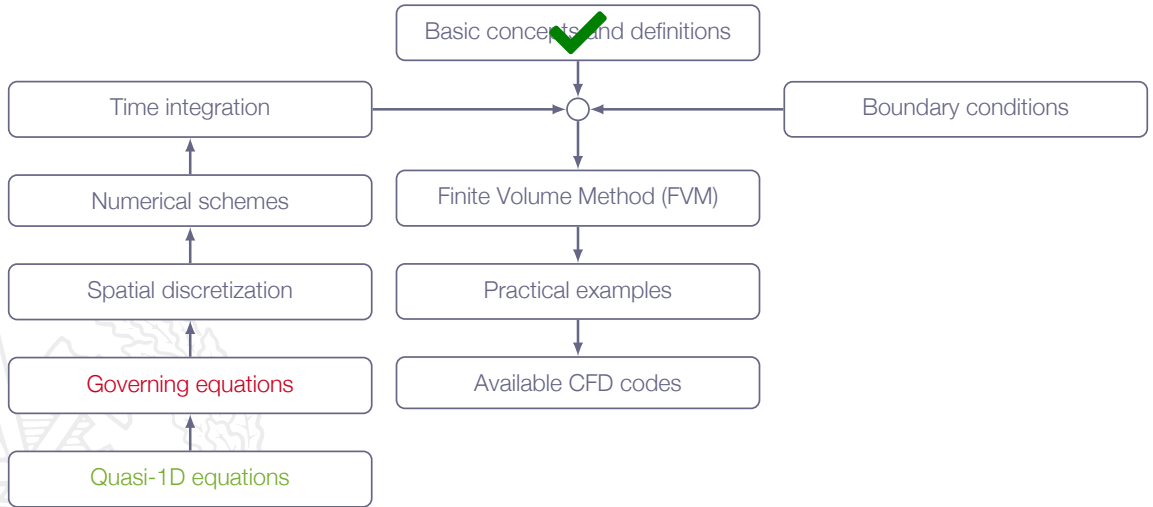
**Explicit Time
Stepping**

**Implicit Time
Stepping**

In general implicit solvers are more efficient than explicit solvers

For high-supersonic flows, explicit solvers may very well outperform implicit solvers

Roadmap - The Time-Marching Technique

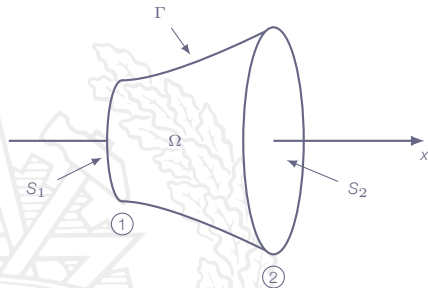


Governing Equations



Introduce **cross-section-averaged flow quantities** $\Rightarrow$
all quantities depend on x only

$$A = A(x), \rho = \rho(x), u = u(x), p = p(x), \dots$$



Ω	control volume
S_1	left boundary (area A_1)
S_2	right boundary (area A_2)
Γ	perimeter boundary

$$\partial\Omega = S_1 \cup \Gamma \cup S_2$$

Governing equations (general form):

$$\frac{d}{dt} \iiint_{\Omega} \rho d\mathcal{V} + \iint_{\partial\Omega} \rho \mathbf{v} \cdot \mathbf{n} dS = 0$$

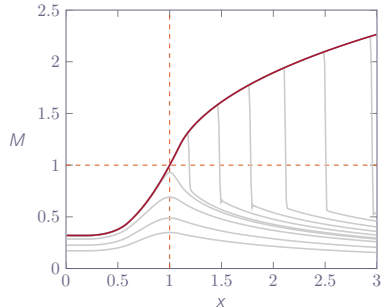
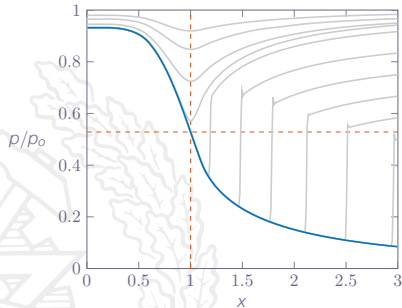
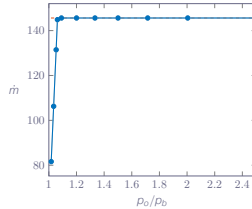
$$\frac{d}{dt} \iiint_{\Omega} \rho u d\mathcal{V} + \iint_{\partial\Omega} [\rho(\mathbf{v} \cdot \mathbf{n})u + p(\mathbf{n} \cdot \mathbf{e}_x)] dS = 0$$

$$\frac{d}{dt} \iiint_{\Omega} \rho e_o d\mathcal{V} + \iint_{\partial\Omega} \rho h_o(\mathbf{v} \cdot \mathbf{n}) dS = 0$$

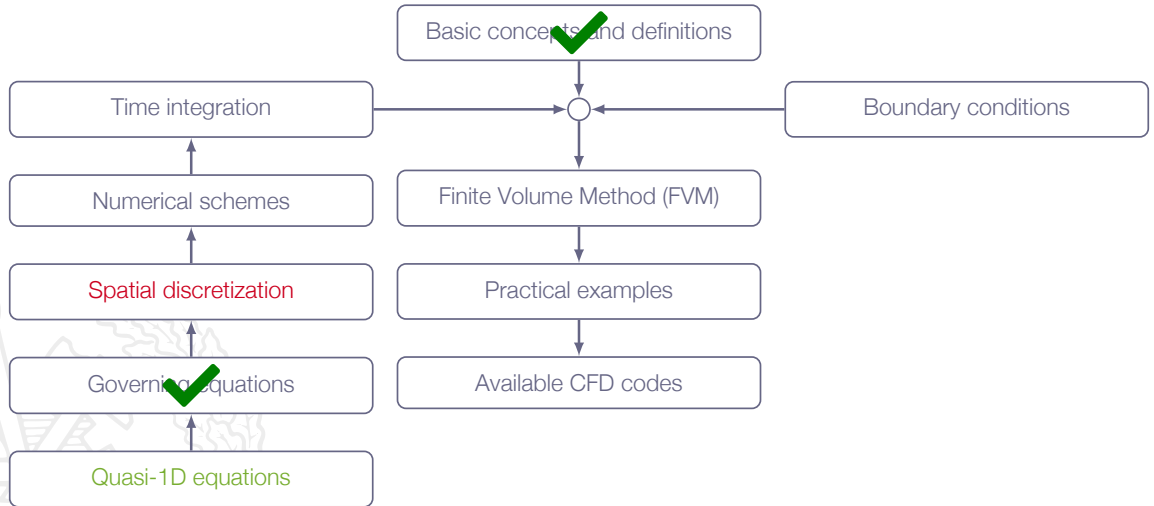


Quasi-One-Dimensional Flow - Example: Nozzle Flow

p_o	1.20 [bar]
p_b	0.50 [bar]
p_o/p_b	11.8
$\dot{m}$	145.6 [kg/s]
M_{max}	2.26



Roadmap - The Time-Marching Technique



Spatial Discretization



Discretization in space and time:

Method of Lines (a very common approach):

1. discretize in space $\Rightarrow$ system of ordinary differential equations (ODEs)
2. discretize in time $\Rightarrow$ time-stepping scheme for system of ODEs

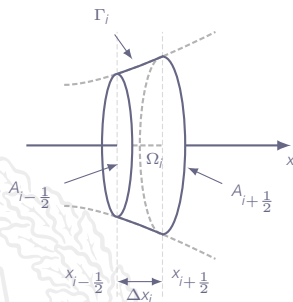
Spatial discretization techniques:

FDM Finite-Difference Method

FVM **Finite-Volume Method**

FEM Finite-Element Method

Let's look at a small tube segment with length Δx



Streamtube with area $A(x)$

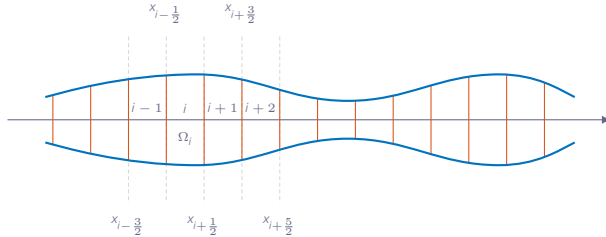
$$A_{i-\frac{1}{2}} = A(x_{i-\frac{1}{2}})$$

$$A_{i+\frac{1}{2}} = A(x_{i+\frac{1}{2}})$$

$$\Delta x_i = x_{i+\frac{1}{2}} - x_{i-\frac{1}{2}}$$

Ω_i - control volume enclosed by $A_{i-\frac{1}{2}}$, $A_{i+\frac{1}{2}}$, and Γ_i

$\Rightarrow$ **spatial discretization**



Integer indices: control volumes or **cells**

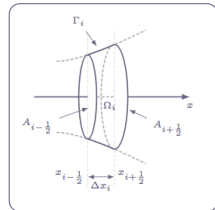
Fractional indices: interfaces between control volumes or **cell faces**

Apply control volume formulations for mass, momentum, energy to control volume Ω_i

Quasi-One-Dimensional Flow - Spatial Discretization

cell-averaged quantity

face-averaged quantity



Conservation of mass:

$$\underbrace{\frac{d}{dt} \iiint_{\Omega_i} \rho d\mathcal{V}}_{VOL_i \frac{d}{dt} \bar{\rho}_i} + \underbrace{\iint_{x_{i-\frac{1}{2}}} \rho \mathbf{v} \cdot \mathbf{n} dS}_{-\overline{(\rho u)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}}} + \underbrace{\iint_{x_{i+\frac{1}{2}}} \rho \mathbf{v} \cdot \mathbf{n} dS}_{\overline{(\rho u)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}}} + \underbrace{\iint_{\Gamma_i} \rho \mathbf{v} \cdot \mathbf{n} dS}_0 = 0$$

where

$$VOL_i = \iiint_{\Omega_i} d\mathcal{V}$$

$$\bar{\rho}_i = \frac{1}{VOL_i} \iiint_{\Omega_i} \rho d\mathcal{V}$$

$$\overline{(\rho u)}_{i-\frac{1}{2}} = \frac{1}{A_{i-\frac{1}{2}}} \iint_{x_{i-\frac{1}{2}}} \rho u dS$$

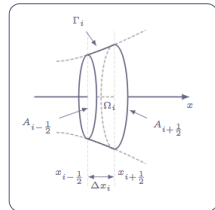
$$\overline{(\rho u)}_{i+\frac{1}{2}} = \frac{1}{A_{i+\frac{1}{2}}} \iint_{x_{i+\frac{1}{2}}} \rho u dS$$

Quasi-One-Dimensional Flow - Spatial Discretization

cell-averaged quantity

face-averaged quantity

source term



Conservation of momentum:

$$\underbrace{\frac{d}{dt} \iiint_{\Omega_i} \rho u d\mathcal{V}}_{VOL_i \frac{d}{dt} \overline{(\rho u)}_i} + \underbrace{\iint_{x_{i-\frac{1}{2}}} [\rho(\mathbf{v} \cdot \mathbf{n})u + p(\mathbf{n} \cdot \mathbf{e}_x)] dS}_{-\overline{(\rho u^2 + p)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}}} +$$

$$+ \underbrace{\iint_{x_{i+\frac{1}{2}}} [\rho(\mathbf{v} \cdot \mathbf{n})u + p(\mathbf{n} \cdot \mathbf{e}_x)] dS}_{\overline{(\rho u^2 + p)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}}} + \underbrace{\iint_{\Gamma_i} [\rho(\mathbf{v} \cdot \mathbf{n})u + p(\mathbf{n} \cdot \mathbf{e}_x)] dS}_{-\iint_{\Gamma_i} p dA} = 0$$

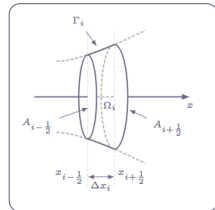
Quasi-One-Dimensional Flow - Spatial Discretization

cell-averaged quantity

face-averaged quantity

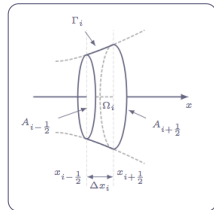
Conservation of energy:

$$\begin{aligned}
 & \underbrace{\frac{d}{dt} \iiint_{\Omega_i} \rho e_o d\mathcal{V}}_{VOL_i \frac{d}{dt} \overline{(\rho e_o)}_i} + \underbrace{\iint_{x_{i-\frac{1}{2}}} \rho h_o (\mathbf{v} \cdot \mathbf{n}) dS}_{-\overline{(\rho u h_o)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}}} \\
 & + \underbrace{\iint_{x_{i+\frac{1}{2}}} \rho h_o (\mathbf{v} \cdot \mathbf{n}) dS}_{\overline{(\rho u h_o)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}}} + \underbrace{\iint_{\Gamma_i} \rho h_o (\mathbf{v} \cdot \mathbf{n}) dS}_0 = 0
 \end{aligned}$$



Quasi-One-Dimensional Flow - Spatial Discretization

Lower order term due to varying stream tube area:



$$\iint_{\Gamma_i} p dA \approx \bar{p}_i \left(A_{i+\frac{1}{2}} - A_{i-\frac{1}{2}} \right)$$

where $\bar{p}_i$ is **calculated from cell-averaged quantities** (DOFs) $\left\{ \bar{\rho}, \overline{(\rho u)}, \overline{(\rho e_o)} \right\}_i$ as

$$\bar{p}_i = (\gamma - 1) \left(\overline{(\rho e_o)}_i - \frac{1}{2} \bar{\rho}_i \bar{u}_i^2 \right), \quad \bar{u}_i = \frac{\overline{(\rho u)}_i}{\bar{\rho}_i}$$

cell-averaged quantity

face-averaged quantity

source term

$$VOL_i \frac{d}{dt} \bar{\rho}_i - \overline{(\rho u)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} + \overline{(\rho u)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} = 0$$

$$VOL_i \frac{d}{dt} \overline{(\rho u)}_i - \overline{(\rho u^2 + p)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} + \overline{(\rho u^2 + p)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} = \bar{p}_i \left(A_{i+\frac{1}{2}} - A_{i-\frac{1}{2}} \right)$$

$$VOL_i \frac{d}{dt} \overline{(\rho e_o)}_i - \overline{(\rho u h_o)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} + \overline{(\rho u h_o)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} = 0$$

Application of these equations to all cells $i \in \{1, 2, \dots, N\}$ of the computational domain results in a system of ODEs

Steps to achieve spatial discretization:

1. Choose primary variables (degrees of freedom)
2. Approximate all other quantities in terms of the primary variables

⇒ **System of ordinary differential equations** (ODEs)

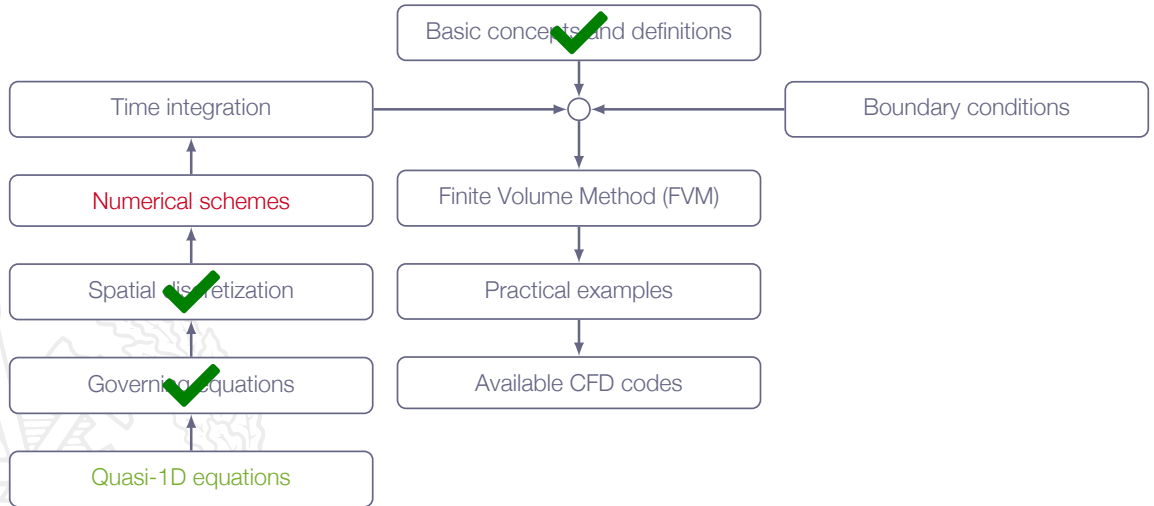
Degrees of freedom:

Choose $\left\{ \bar{\rho}, \overline{(\rho u)}, \overline{(\rho e_o)} \right\}_i$ in all control volumes $\Omega_i, i \in \{1, 2, \dots, N\}$ as degrees of freedom, or primary variables

Note that these are **cell-averaged quantities**

What about the face values?

Roadmap - The Time-Marching Technique



Numerical Schemes



$$\left\{ \begin{array}{c} \overline{(\rho u)} \\ \overline{(\rho u^2 + p)} \\ \overline{(\rho u h_o)} \end{array} \right\}_{i+\frac{1}{2}} = f \left(\left\{ \begin{array}{c} \bar{\rho} \\ \overline{(\rho u)} \\ \overline{(\rho e_o)} \end{array} \right\}_i, \left\{ \begin{array}{c} \bar{\rho} \\ \overline{(\rho u)} \\ \overline{(\rho e_o)} \end{array} \right\}_{i+1}, \dots \right)$$

cell face values

cell-averaged values

Simple example:

$$\overline{(\rho u)}_{i+\frac{1}{2}} \approx \frac{1}{2} \left[\overline{(\rho u)}_i + \overline{(\rho u)}_{i+1} \right]$$

More complex approximations usually needed

High-order schemes:

- increased accuracy

- more cell values involved (*wider flux molecule*)

- boundary conditions more difficult to implement

Optimized numerical dissipation:

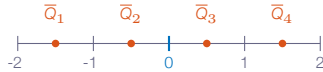
- upwind type of flux scheme

Shock handling:

- non-linear treatment needed (e.g. TVD schemes)

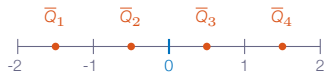
- artificial damping

Flux Term Approximation



$$Q(x) = A + Bx + Cx^2 + Dx^3$$

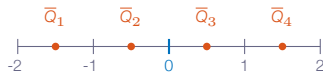
Assume constant area: $A(x) = 1.0$



$$\bar{Q}_1 = \frac{1}{VOL_1} \int_{-2}^{-1} Q(x) dx$$

$$VOL_1 = A_1 \Delta x_1 = \{A_1 = 1.0, \Delta x_1 = 1.0\} = 1.0$$

$$\Rightarrow \bar{Q}_1 = \int_{-2}^{-1} Q(x) dx$$

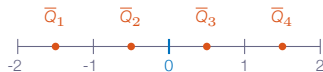


$$\bar{Q}_1 = \int_{-2}^{-1} Q(x) dx = \left[Ax + \frac{1}{2}Bx^2 + \frac{1}{3}Cx^3 + \frac{1}{4}Dx^4 \right]_{-2}^{-1}$$

$$\bar{Q}_2 = \int_{-1}^0 Q(x) dx = \left[Ax + \frac{1}{2}Bx^2 + \frac{1}{3}Cx^3 + \frac{1}{4}Dx^4 \right]_{-1}^0$$

$$\bar{Q}_3 = \int_0^1 Q(x) dx = \left[Ax + \frac{1}{2}Bx^2 + \frac{1}{3}Cx^3 + \frac{1}{4}Dx^4 \right]_0^1$$

$$\bar{Q}_4 = \int_1^2 Q(x) dx = \left[Ax + \frac{1}{2}Bx^2 + \frac{1}{3}Cx^3 + \frac{1}{4}Dx^4 \right]_1^2$$

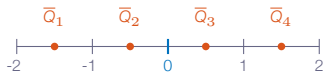


$$\bar{Q}_1 = A - \frac{3}{2}B + \frac{7}{3}C - \frac{15}{4}D$$

$$\bar{Q}_2 = A - \frac{1}{2}B + \frac{1}{3}C - \frac{1}{4}D$$

$$\bar{Q}_3 = A + \frac{1}{2}B + \frac{1}{3}C + \frac{1}{4}D$$

$$\bar{Q}_4 = A + \frac{3}{2}B + \frac{7}{3}C + \frac{15}{4}D$$

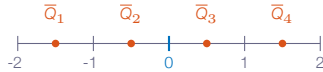


$$A = \frac{1}{12} \left[-\bar{Q}_1 + 7\bar{Q}_2 + 7\bar{Q}_3 - \bar{Q}_4 \right]$$

$$B = \frac{1}{12} \left[\bar{Q}_1 - 15\bar{Q}_2 + 15\bar{Q}_3 - \bar{Q}_4 \right]$$

$$C = \frac{1}{4} \left[\bar{Q}_1 - \bar{Q}_2 - \bar{Q}_3 + \bar{Q}_4 \right]$$

$$D = \frac{1}{6} \left[-\bar{Q}_1 + 3\bar{Q}_2 - 3\bar{Q}_3 + \bar{Q}_4 \right]$$

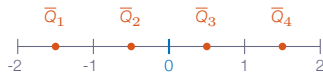


$$Q_0 = Q(0) + \delta Q'''(0) \Rightarrow Q_0 = A + 6\delta D$$

$\delta = 0 \Rightarrow$ fourth-order central scheme

$\delta = 1/12 \Rightarrow$ third-order upwind scheme

$\delta = 1/96 \Rightarrow$ third-order low-dissipation upwind scheme



$$Q_0 = A + 6\delta D = \{\delta = 1/12\} = -\frac{1}{6}\bar{Q}_1 + \frac{5}{6}\bar{Q}_2 + \frac{1}{3}\bar{Q}_3$$

$$Q_{0_{left}} = -\frac{1}{6}\bar{Q}_1 + \frac{5}{6}\bar{Q}_2 + \frac{1}{3}\bar{Q}_3$$

$$Q_{0_{right}} = -\frac{1}{6}\bar{Q}_4 + \frac{5}{6}\bar{Q}_3 + \frac{1}{3}\bar{Q}_2$$

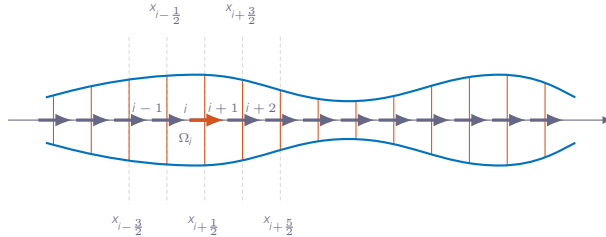
method of characteristics used in order to decide whether left- or right-upwinded flow quantities should be used

High-order numerical schemes:

low numerical dissipation (smearing due to amplitudes errors)

low dispersion errors (wiggles due to phase errors)





mass conservation:

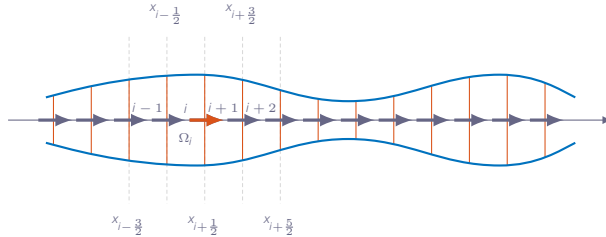
cell (i):

$$VOL_i \frac{d}{dt} \bar{\rho}_i + \overline{(\rho u)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} - \overline{(\rho u)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} = 0$$

cell ($i+1$):

$$VOL_{i+1} \frac{d}{dt} \bar{\rho}_{i+1} + \overline{(\rho u)}_{i+\frac{3}{2}} A_{i+\frac{3}{2}} - \overline{(\rho u)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} = 0$$

(similarly for momentum and energy conservation)



mass conservation:

cell (i):

$$VOL_i \frac{d}{dt} \bar{\rho}_i + \overline{(\rho u)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} - \overline{(\rho u)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} = 0$$

cell ($i+1$):

$$VOL_{i+1} \frac{d}{dt} \bar{\rho}_{i+1} + \overline{(\rho u)}_{i+\frac{3}{2}} A_{i+\frac{3}{2}} - \overline{(\rho u)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} = 0$$

(similarly for momentum and energy conservation)

Conservative scheme

"The flux leaving one control volume equals the flux entering neighbouring control volume"

Conservation of for mass, momentum and energy is crucial for the correct prediction of shocks*

* correct prediction of shocks:
strength
position
velocity

Jameson shock detector:

$$\nu_{i+\frac{1}{2}} = \max \{ \nu_i, \nu_{i+1} \}$$

where ν_i is a scaled pressure derivative

$$\nu_i = \frac{|p_{i+1} - 2p_i + p_{i-1}|}{p_{i+1} + 2p_i + p_{i-1}}$$

For a smooth pressure field $\nu \mathcal{O}(\Delta x^2)$ and near a shock $\nu \mathcal{O}(1)$

Artificial damping term (α is a user-defined constant):

$$\alpha (|u| + c)_{i+\frac{1}{2}} \nu_{i+\frac{1}{2}} A_{i+\frac{1}{2}} (Q_{i+1} - Q_i)$$

Jameson-type detector:

$$\nu_{i+\frac{1}{2}} = \max \{ \nu_i, \nu_{i+1} \}$$

where ν_i is a scaled density derivative

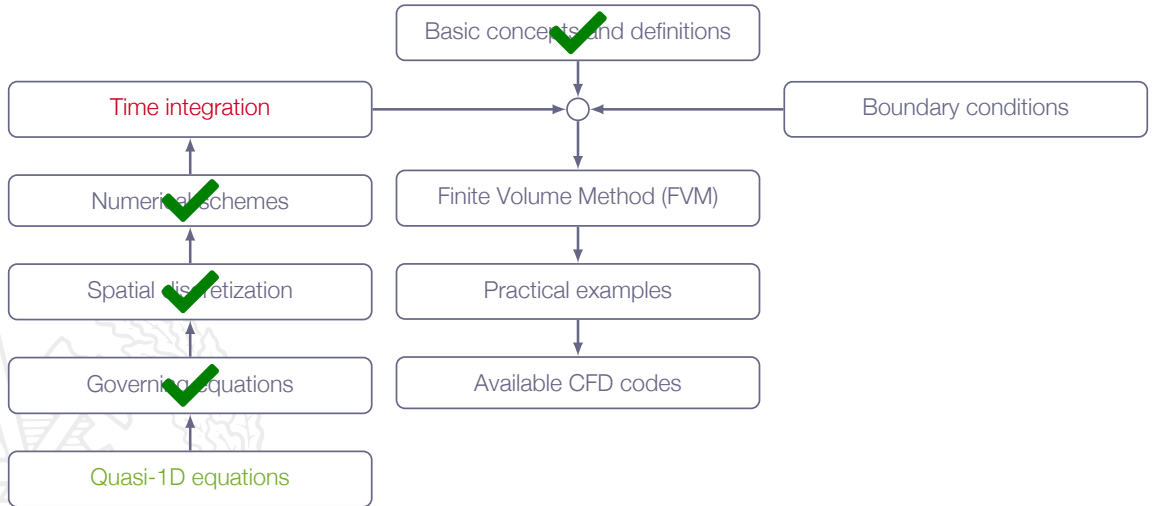
$$\nu_i = \frac{|\rho_{i+1} - 2\rho_i + \rho_{i-1}|}{\rho_{i+1} + 2\rho_i + \rho_{i-1}}$$

For a smooth density field $\nu \mathcal{O}(\Delta x^2)$ and near a density discontinuity $\nu \mathcal{O}(1)$

Artificial damping term (β is a user-defined constant):

$$\beta |u|_{i+\frac{1}{2}} \nu_{i+\frac{1}{2}} A_{i+\frac{1}{2}} (Q_{i+1} - Q_i)$$

Roadmap - The Time-Marching Technique



Time Stepping



cell-averaged quantity

face-averaged quantity

source term

$$VOL_i \frac{d}{dt} \bar{\rho}_i - \overline{(\rho u)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} + \overline{(\rho u)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} = 0$$

$$VOL_i \frac{d}{dt} \overline{(\rho u)}_i - \overline{(\rho u^2 + p)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} + \overline{(\rho u^2 + p)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} = \bar{p}_i (A_{i+\frac{1}{2}} - A_{i-\frac{1}{2}})$$

$$VOL_i \frac{d}{dt} \overline{(\rho e_o)}_i - \overline{(\rho u h_o)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} + \overline{(\rho u h_o)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} = 0$$

Application of these equations to all cells $i \in \{1, 2, \dots, N\}$ of the computational domain results in a system of ODEs

cell-averaged quantity

face-averaged quantity

source term

$$VOL_i \frac{d}{dt} \bar{\rho}_i = \overline{(\rho u)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} - \overline{(\rho u)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}}$$

$$VOL_i \frac{d}{dt} \overline{(\rho u)}_i = \overline{(\rho u^2 + p)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} - \overline{(\rho u^2 + p)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} + \bar{p}_i \left(A_{i+\frac{1}{2}} - A_{i-\frac{1}{2}} \right)$$

$$VOL_i \frac{d}{dt} \overline{(\rho e_o)}_i = \overline{(\rho u h_o)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} - \overline{(\rho u h_o)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}}$$

cell-averaged quantity

face-averaged quantity

source term

$$\frac{d}{dt} \bar{\rho}_i = \frac{1}{VOL_i} \left[\overline{(\rho u)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} - \overline{(\rho u)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} \right]$$

$$\frac{d}{dt} \overline{(\rho u)}_i = \frac{1}{VOL_i} \left[\overline{(\rho u^2 + p)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} - \overline{(\rho u^2 + p)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} + \bar{p}_i \left(A_{i+\frac{1}{2}} - A_{i-\frac{1}{2}} \right) \right]$$

$$\frac{d}{dt} \overline{(\rho e_o)}_i = \frac{1}{VOL_i} \left[\overline{(\rho u h_o)}_{i-\frac{1}{2}} A_{i-\frac{1}{2}} - \overline{(\rho u h_o)}_{i+\frac{1}{2}} A_{i+\frac{1}{2}} \right]$$

$$\frac{d}{dt} \bar{\mathbf{Q}}_i = \mathbf{F}(\bar{\mathbf{Q}}_i) \text{ where } \bar{\mathbf{Q}}_i = [\bar{\rho}, \bar{\rho u}, \overline{\rho e_o}]_i, i \in \{1 : NCells\}$$

The system of ODEs obtained from the spatial discretization in vector notation

$$\frac{d}{dt}\mathbf{Q} = \mathbf{F}(\mathbf{Q})$$

$\mathbf{Q}$ is a vector containing all DOFs in all cells

$\mathbf{F}(\mathbf{Q})$ is the **time derivative** of $\mathbf{Q}$ resulting from above mentioned **flux approximations** - *non-linear vector-valued function*

Three-stage Runge-Kutta - *one example of many:*

Explicit time-marching scheme

Second-order accurate



$$\frac{d}{dt}\mathbf{Q} = \mathbf{F}(\mathbf{Q})$$

Let $\mathbf{Q}^n = \mathbf{Q}(t_n)$ and $\mathbf{Q}^{n+1} = \mathbf{Q}(t_{n+1})$

t_n is the current time level and t_{n+1} is the next time level

$\Delta t = t_{n+1} - t_n$ is the solver time step

Algorithm:

1. $\mathbf{Q}^* = \mathbf{Q}^n + \Delta t \mathbf{F}(\mathbf{Q}^n)$
2. $\mathbf{Q}^{**} = \mathbf{Q}^n + \frac{1}{2} \Delta t \mathbf{F}(\mathbf{Q}^n) + \frac{1}{2} \Delta t \mathbf{F}(\mathbf{Q}^*)$
3. $\mathbf{Q}^{n+1} = \mathbf{Q}^n + \frac{1}{2} \Delta t \mathbf{F}(\mathbf{Q}^n) + \frac{1}{2} \Delta t \mathbf{F}(\mathbf{Q}^{**})$

```
1 void RungeKutta::fwd(Domain *dom){
2     G3DCopy(dom->cons, cons0);
3
4     /* Runge-Kutta step 1 */
5
6     dom->update();
7     if(!G3DMode::constdt){LocalTimeStep(dom);}
8     dcons->evaluate(dom);
9     G3DWXPY(dom->cons, 1.0, dcons, cons0);
10    G3DAXPY(cons0, 0.5, 0.5, dom->cons);
11
12    /* Runge-Kutta step 2 */
13
14    dom->update();
15    dcons->evaluate(dom);
16    G3DWXPY(dom->cons, 0.5, dcons, cons0);
17
18    /* Runge-Kutta step 3 */
19
20    dom->update();
21    dcons->evaluate(dom);
22    G3DWXPY(dom->cons, 0.5, dcons, cons0);
23 }
```

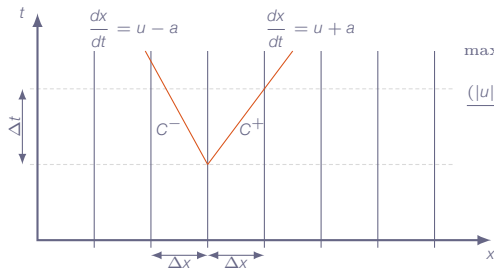
Properties of explicit time-stepping schemes:

- + **Easy to implement** in computer codes
- + **Efficient execution** on most computers
- + Easy to adapt for **parallel execution** on distributed memory systems (e.g. Linux clusters)
- **Time step limitation** (CFL number)
- Convergence to steady-state **often slow** (there are, however, some remedies for this)

Courant-Friedrich-Levy (**CFL**) number - *one-dimensional case*:

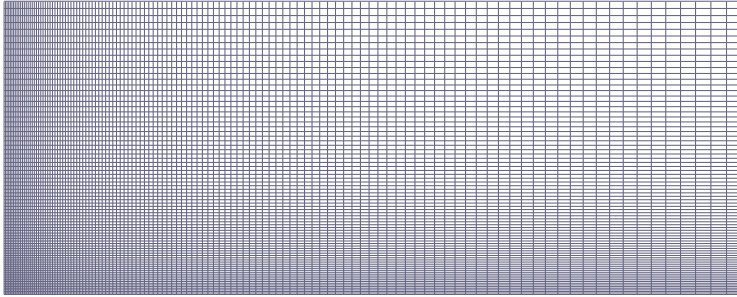
$$CFL_i = \frac{\Delta t(|u_i| + a_i)}{\Delta x_i} \leq 1$$

Interpretation: The fastest characteristic (C^+ or C^-) must not travel longer than Δx during one time step



$$\max(|u - a|, |u + a|)\Delta t = (|u| + a)\Delta t \leq \Delta x \Rightarrow$$

$$\frac{(|u| + a)\Delta t}{\Delta x} = CFL \leq 1$$



Steady-state problems:

- local time stepping

- each cell has an individual time step

- Δt_i maximum allowed value based on CFL criteria

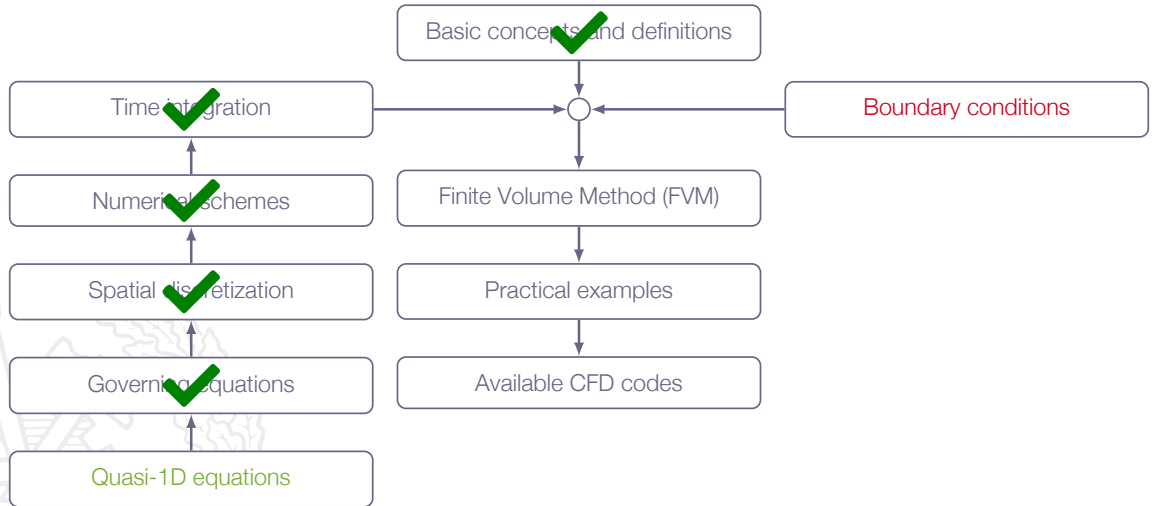
Unsteady problems:

- time accurate

- all cells have the same time step

$$\Delta t_i = \min \{ \Delta t_1, \dots, \Delta t_N \}$$

Roadmap - The Time-Marching Technique



Boundary Conditions



Boundary conditions are very important for numerical simulation of compressible flows

Main reason: both **flow** and **acoustics** involved!

Example 1:

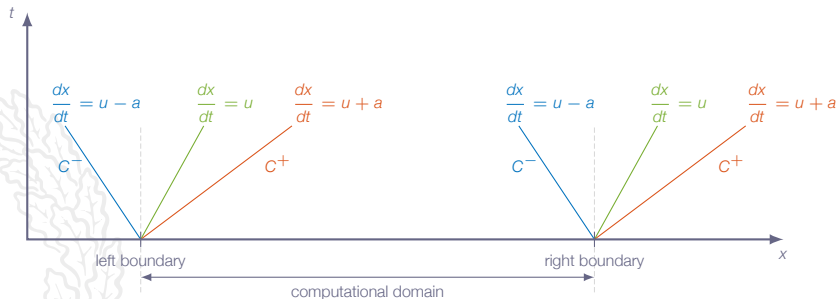
Finite-volume CFD code for Quasi-1D compressible flow (Time-marching procedure)

What boundary conditions should be applied at the left and right ends?



three characteristics:

1. C^+
2. C^-
3. advection



C^+ and C^- characteristics describe the transport of **isentropic pressure waves** (often referred to as **acoustics**)

The advection characteristic simply describes the **transport** of certain quantities **with the fluid itself** (for example **entropy**)

In one space dimension and time, these three characteristics, together with the quantities that are known to be constant along them, give a **complete description** of the time evolution of the flow

We can use the characteristics as a guide to tell us what information that should be specified at the boundaries

we have three PDEs, and are solving for three unknowns

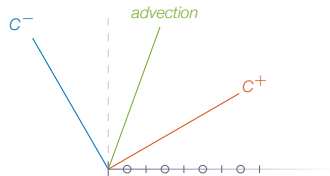
Subsonic inflow: $0 < u < a$

$$u - a < 0$$

$$u > 0$$

$$u + a > 0$$

one outgoing characteristic
two ingoing characteristics



Two variables should be **specified** at the boundary

The third variable must be left free

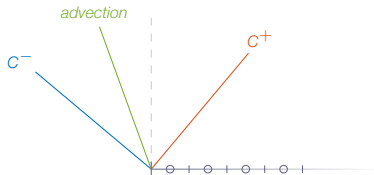
we have three PDEs, and are solving for three unknowns

Subsonic outflow: $-a < u < 0$

$$u - a < 0$$

$$u < 0$$

$$u + a > 0$$



two outgoing characteristics
one ingoing characteristic

One variable should be **specified** at the boundary

The second and third variables must be left free

we have three PDEs, and are solving for three unknowns

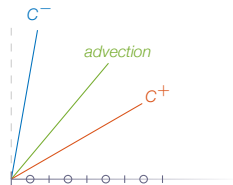
Supersonic inflow: $u > a$

$$u - a > 0$$

$$u > 0$$

$$u + a > 0$$

no outgoing characteristics
three ingoing characteristics



All three variables should be **specified** at the boundary

No variables must be left free

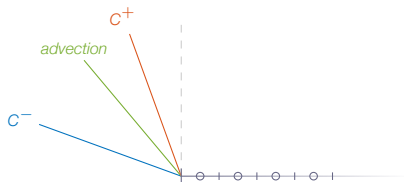
we have three PDEs, and are solving for three unknowns

Supersonic outflow: $u < -a$

$$u - a < 0$$

$$u < 0$$

$$u + a < 0$$



three outgoing characteristics
no ingoing characteristics

No variables should be **specified** at the boundary

All variables must be left free

we have three PDEs, and are solving for three unknowns

Subsonic inflow: $-a < u < 0$

$$u - a < 0$$

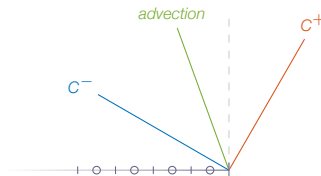
$$u < 0$$

$$u + a > 0$$

two ingoing characteristics
one outgoing characteristic

Two variables should be **specified** at the boundary

The third variables must be left free



we have three PDEs, and are solving for three unknowns

Subsonic outflow: $0 < u < a$

$$u - a < 0$$

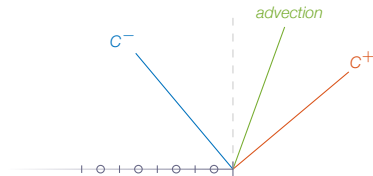
$$u > 0$$

$$u + a > 0$$

one ingoing characteristic
two outgoing characteristics

One variable should be **specified** at the boundary

The second and third variables must be left free



we have three PDEs, and are solving for three unknowns

Supersonic inflow: $u < -a$

$$u - a < 0$$

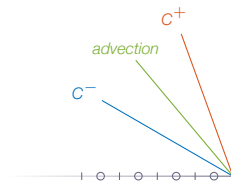
$$u < 0$$

$$u + a < 0$$

three ingoing characteristics
no outgoing characteristics

All three variables should be **specified** at the boundary

No variables must be left free



we have three PDEs, and are solving for three unknowns

Supersonic outflow: $u > a$

$$u - a > 0$$

$$u > 0$$

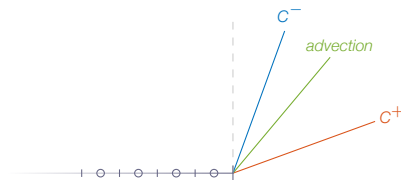
$$u + a > 0$$

no ingoing characteristics

three outgoing characteristics

No variables should be **specified** at the boundary

All three variables must be left free



1D Boundary Conditions (Summary)

Characteristic		1D subsonic inflow (left)	1D subsonic inflow (right)
advection	$\mathbf{v} \cdot \mathbf{n}$	$(u, 0, 0) \cdot (-1, 0, 0) = -u < 0$	$(-u, 0, 0) \cdot (1, 0, 0) = -u < 0$
C^-	$\mathbf{v} \cdot \mathbf{n} - a$	$-u - a < 0$	$-u - a < 0$
C^+	$\mathbf{v} \cdot \mathbf{n} + a$	$-u + a > 0$	$-u + a > 0$
Characteristic		1D subsonic outflow (left)	1D subsonic outflow (right)
advection	$\mathbf{v} \cdot \mathbf{n}$	$(-u, 0, 0) \cdot (-1, 0, 0) = u > 0$	$(u, 0, 0) \cdot (1, 0, 0) = u > 0$
C^-	$\mathbf{v} \cdot \mathbf{n} - a$	$u - a < 0$	$u - a < 0$
C^+	$\mathbf{v} \cdot \mathbf{n} + a$	$u + a > 0$	$u + a > 0$
Characteristic		1D supersonic inflow (left)	1D supersonic inflow (right)
advection	$\mathbf{v} \cdot \mathbf{n}$	$(u, 0, 0) \cdot (-1, 0, 0) = -u < 0$	$(-u, 0, 0) \cdot (1, 0, 0) = -u < 0$
C^-	$\mathbf{v} \cdot \mathbf{n} - a$	$-u - a < 0$	$-u - a < 0$
C^+	$\mathbf{v} \cdot \mathbf{n} + a$	$-u + a < 0$	$-u + a < 0$
Characteristic		1D supersonic outflow (left)	1D supersonic outflow (right)
advection	$\mathbf{v} \cdot \mathbf{n}$	$(-u, 0, 0) \cdot (-1, 0, 0) = u > 0$	$(u, 0, 0) \cdot (1, 0, 0) = u > 0$
C^-	$\mathbf{v} \cdot \mathbf{n} - a$	$u - a > 0$	$u - a > 0$
C^+	$\mathbf{v} \cdot \mathbf{n} + a$	$u + a > 0$	$u + a > 0$

Subsonic Inflow (Left Boundary) - Example

Subsonic inflow: we should specify two variables

Alt	specified variable 1	specified variable 2	well-posed	non-reflective
1	p_o	T_o	X	
2	ρu	T_o	X	
3	s	J^+	X	X

well posed:

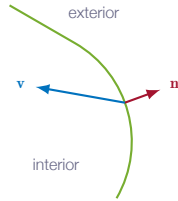
1. the problem has a solution
2. the solution is unique
3. the solution's behaviour changes continuously with initial conditions

Subsonic outflow: we should specify one variable

Alt	specified variable	well-posed	non-reflective
1	p	X	
2	ρu	X	
3	J^+	X	X



Subsonic Inflow 2D/3D



n unit normal vector
v fluid velocity at boundary

Subsonic inflow

Assumption:

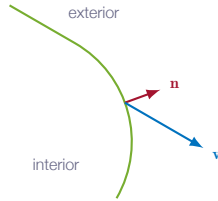
$$-a < \mathbf{v} \cdot \mathbf{n} < 0$$

Four ingoing characteristics

One outgoing characteristic

Specify four variables at the boundary:
 p_o , T_o , and flow direction (two angles)

Subsonic Outflow 2D/3D



n unit normal vector
v fluid velocity at boundary

Subsonic outflow

Assumption:

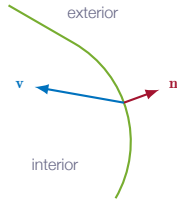
$$0 < \mathbf{v} \cdot \mathbf{n} < a$$

One ingoing characteristics

Four outgoing characteristic

Specify one variables at the boundary:
static pressure

Supersonic Inflow 2D/3D



n unit normal vector
v fluid velocity at boundary

Supersonic inflow

Assumption:

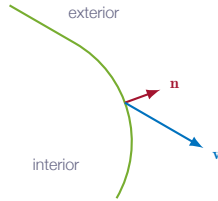
$$\mathbf{v} \cdot \mathbf{n} < -a$$

Five ingoing characteristics

No outgoing characteristics

Specify five variables at the boundary:
solver variables

Supersonic Outflow 2D/3D



$\mathbf{n}$ unit normal vector
 $\mathbf{v}$ fluid velocity at boundary

Supersonic outflow

Assumption:

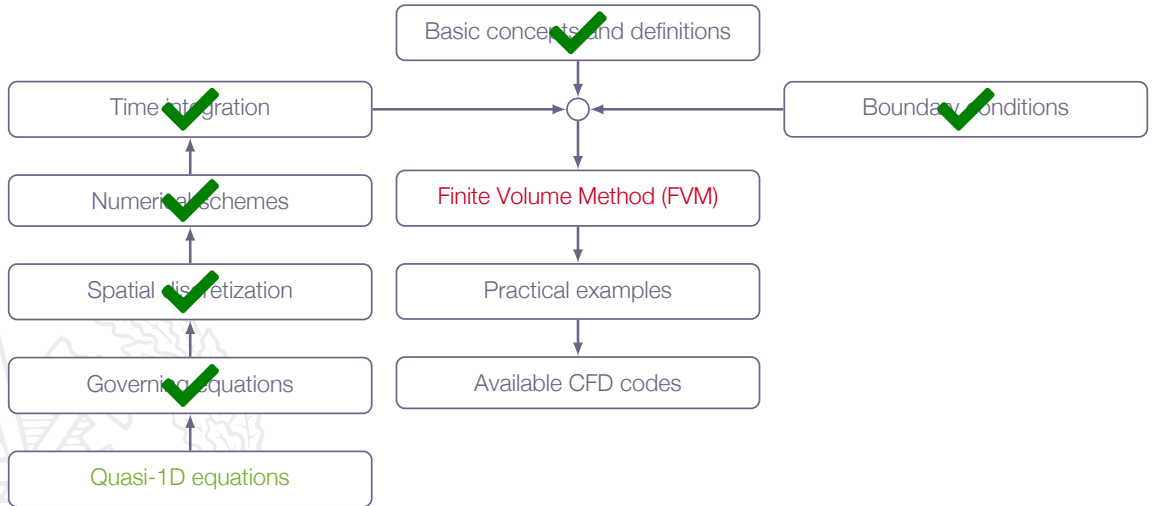
$$\mathbf{v} \cdot \mathbf{n} > a$$

No ingoing characteristics

Five outgoing characteristics

No variables specified at the boundary

Roadmap - The Time-Marching Technique



The described numerical approach can be categorized as

Density-based

Fully coupled

Structured

Explicit

with the following features

**High-order
convective scheme**

**Shock handling
(artificial damping)**

Spatial discretization:

Control volume formulations of conservation equations are applied to the cells of the discretized domain

Cell-averaged flow quantities $(\bar{\rho}, \bar{\rho u}, \bar{\rho e_o})$ are chosen as degrees of freedom

Flux terms are **approximated** in terms of the chosen degrees of freedom
high-order, upwind type of flux approximation is used for optimum results

A **fully conservative** scheme is obtained
the flux leaving one cell is identical to the flux entering the neighboring cell

The result of the spatial discretization is a system of ODEs

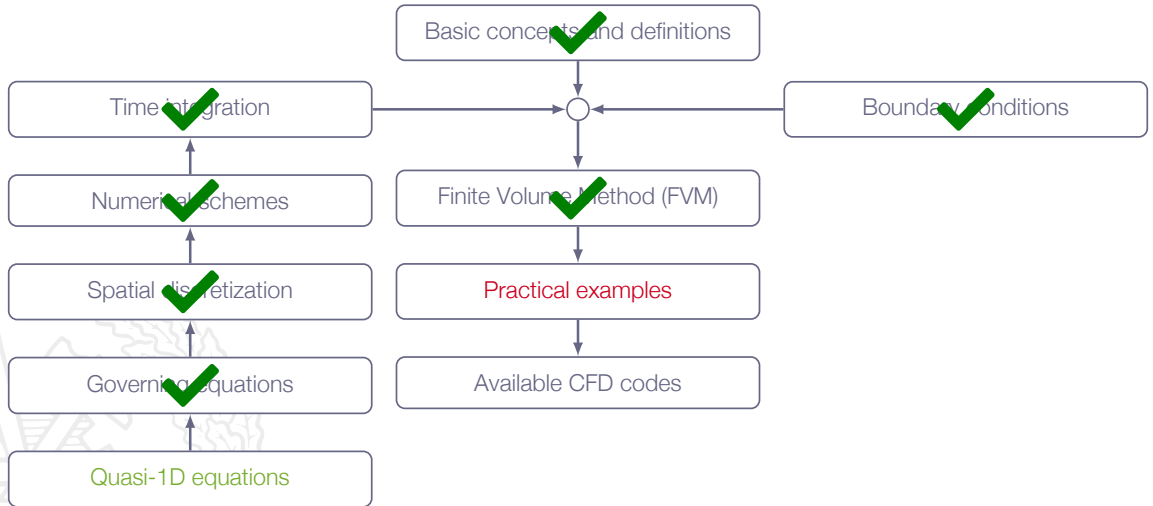
Time marching:

Three-stage, second-order accurate Runge-Kutta scheme

Explicit time-stepping

Time step length **limited by the CFL condition** ($CFL \leq 1$)

Roadmap - The Time-Marching Technique



Practical Examples: Grid Resolution and Numerical Schemes



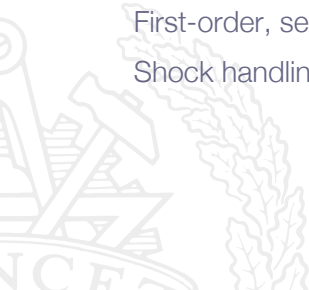
Code: [G3D::Flow](#) (Chalmers in-house CFD code)

Finite-Volume Method

Three-stage, **second-order** accurate **Runge-Kutta** time stepping

First-order, second-order, and **third-order** characteristic upwinding scheme

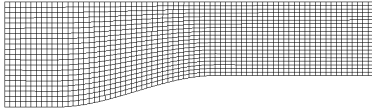
Shock handling: TVD and artificial diffusion based on Jameson shock detection



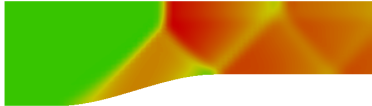
Grid Resolution: Compression Ramp

coarse mesh

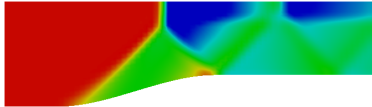
71×21



density

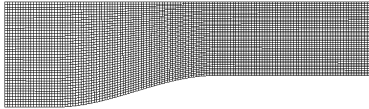


Mach number

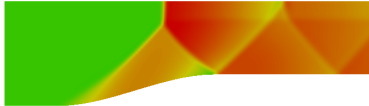


medium mesh

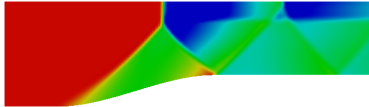
141×41



density

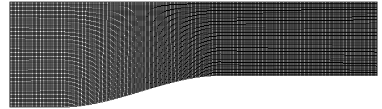


Mach number



fine mesh

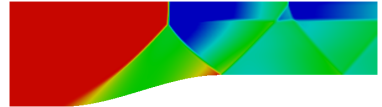
281×81



density



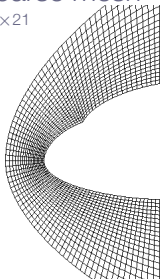
Mach number



Grid Resolution: Space Shuttle

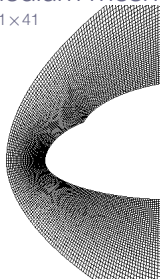
coarse mesh

81×21



medium mesh

161×41

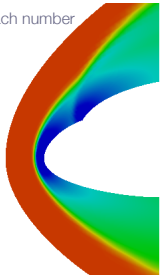


fine mesh

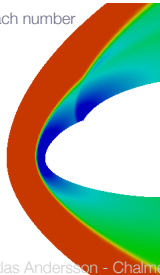
321×81



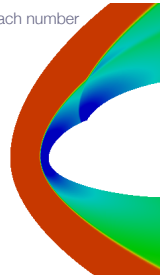
Mach number



Mach number



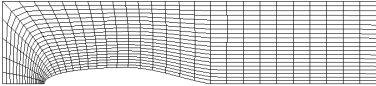
Mach number



Grid Resolution: Axi-symmetric Slender Body

coarse mesh

31×21



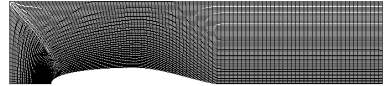
medium mesh

61×41

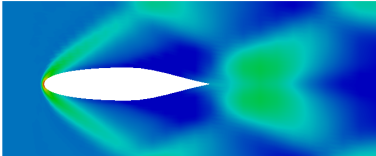


fine mesh

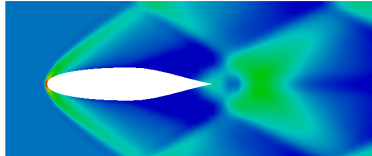
121×81



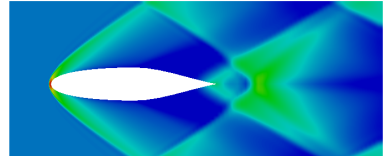
density



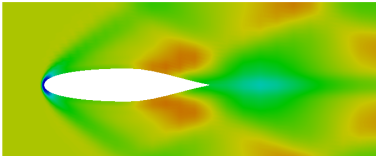
density



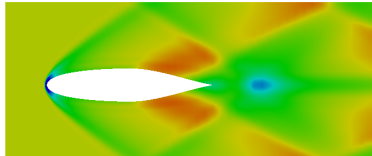
density



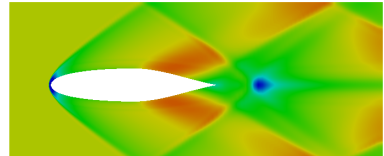
Mach number



Mach number



Mach number



first-order upwind

density



second-order upwind

density

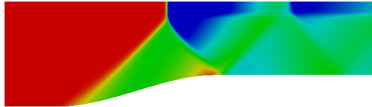


third-order upwind

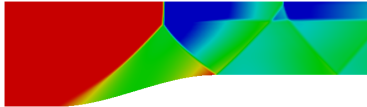
density



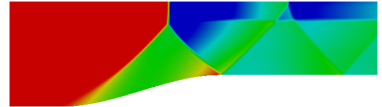
Mach number



Mach number

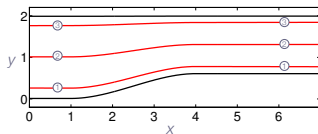


Mach number

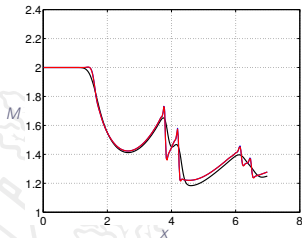


Artificial Numerical Damping: Compression Ramp

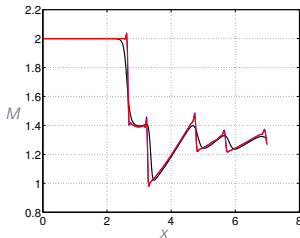
Low artificial numerical damping



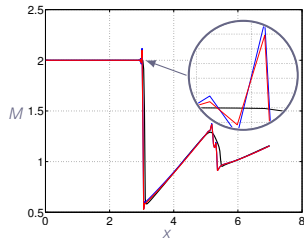
Mach number along line 1



Mach number along line 2



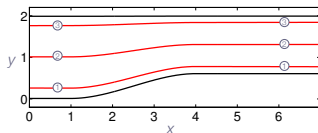
Mach number along line 3



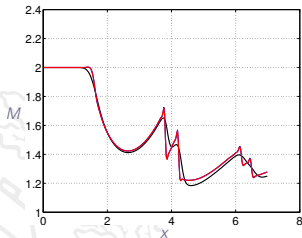
first-order upwind scheme
second-order upwind scheme
third-order upwind scheme

Artificial Numerical Damping: Compression Ramp

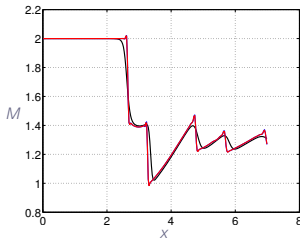
High artificial numerical damping



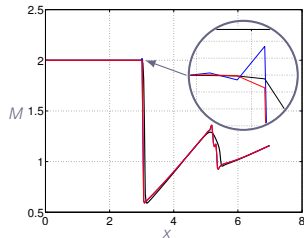
Mach number along line 1



Mach number along line 2

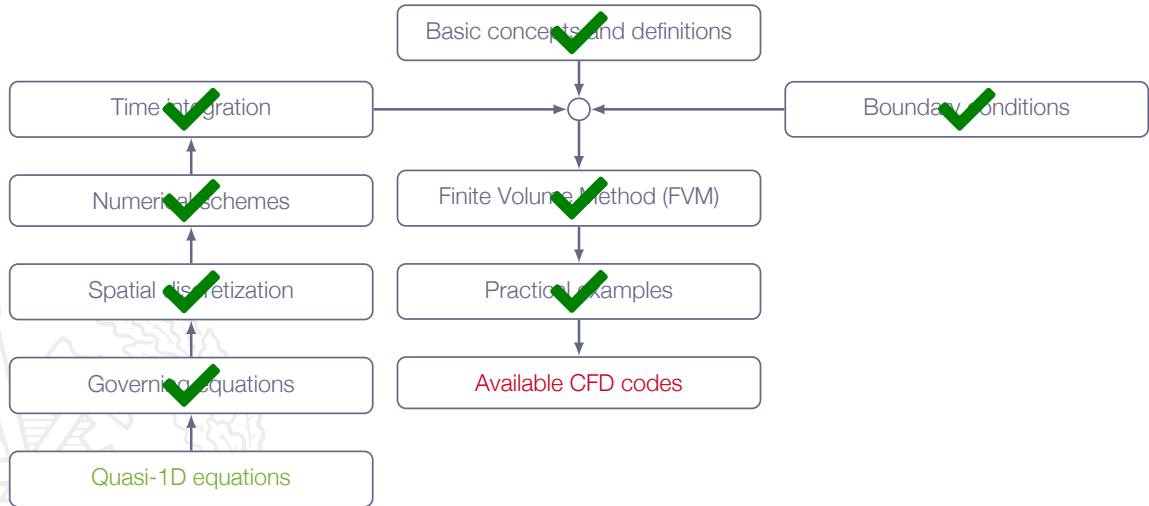


Mach number along line 3



first-order upwind scheme
second-order upwind scheme
third-order upwind scheme

Roadmap - The Time-Marching Technique



Available CFD Codes



CFD Codes

List of free and commercial CFD codes:

<http://www.cfd-online.com/Wiki/Codes>

Free codes are in general unsupported and poorly documented

Commercial codes are often claimed to be suitable for all types of flows

The reality is that the user must make sure of this!



CFD Codes - General Guidelines

Simulation of high-speed and/or unsteady compressible flows:

Use correct solver options

otherwise you may obtain completely wrong solution!

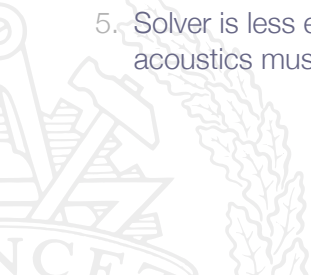
1. coupled solver
2. equation of state
3. energy equation included

Use a high-quality grid

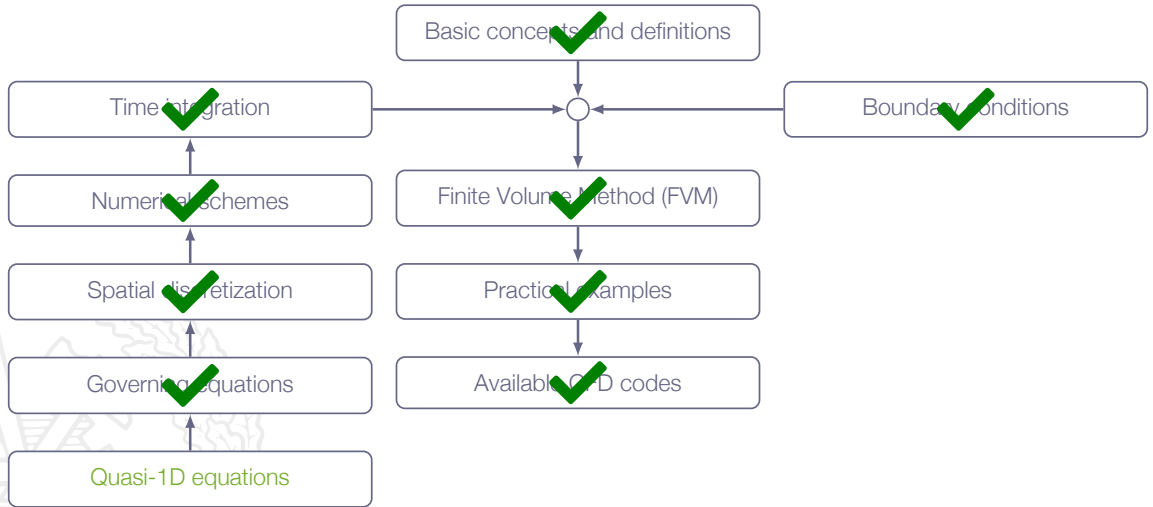
**a poor grid will either not give you any solution at all (no convergence)
or at best a very inaccurate solution!**

ANSYS-FLUENT[®]/STAR-CCM+[®] - Typical Experiences

1. Very robust solvers - will almost always give you a solution
2. Accuracy of solution depends a lot on **grid quality**
3. **Shocks** are generally **smeared** more than in specialized codes
4. Solver is generally very **efficient** for **steady-state** problems
5. Solver is less efficient for truly unsteady problems, where both flow and acoustics must be resolved accurately



Roadmap - The Time-Marching Technique



THE #1 PROGRAMMER EXCUSE
FOR LEGITIMATELY SLACKING OFF:

"MY CODE'S COMPILING."

